No nonlocal box is universal

Frédéric Dupuis* Nicolas Gisin[†] Avinatan Hasidim[§] André Allan Méthot[†] Haran Pilpel[‡]

- * Département d'informatique et de recherche opérationnelle Université de Montréal, Montréal, Québec, CANADA dupuisf@iro.umontreal.ca
- † Group of Applied Physics, Université de Genève, Genève, SWITZERLAND nicolas.gisin, andre.methot@physics.unige.ch
- § Department of Computer Science, Hebrew University, Jerusalem, ISRAEL avinatanh@gmail.com
- [‡] Einstein Institute of Mathematics, Hebrew University, Jerusalem, ISRAEL haranp@math.huji.ac.il

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Abstract

We show that standard nonlocal boxes, also known as Popescu-Rohrlich machines, are not sufficient to simulate any nonlocal correlations that do not allow signalling. This was known in the multipartite scenario, but we extend the result to the bipartite case. We then generalize this result further by showing that *no* finite set containing any finite-output-alphabet nonlocal boxes can be a universal set for nonlocality.

1 Introduction

Nonlocality refers to a multi-party process that, while it does not allow for communication, would classically necessitate communication for the different parties to perform. One classic example is the following "nonlocal box" (see Figure 1): imagine that two parties, henceforth referred to as Alice and Bob, have a black box into which they can each enter one bit of their choice and the box gives each of them a random bit such that the exclusive-or of the output bits is equal to the AND of the input bits [1]. If one attempts to implement this box without communication in a classical world, it is easy to show that it is impossible to succeed more than 75% of the time [2]. On the other hand, if the output bits are always uniformly distributed for all inputs, then it is clear that this box does not permit communication between Alice and Bob, since the local probability distribution does not depend on the parties inputs.

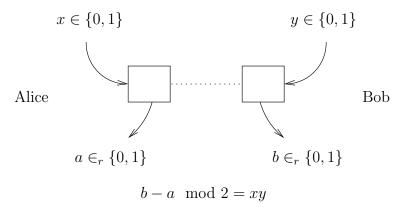


Figure 1: The standard nonlocal box. The notation $a \in_r \{0,1\}$ means that a is uniformly distributed over $\{0,1\}$.

We shall call this box the mod2NLB for reasons that will become obvious later. A general nonlocal box is a device such that: given inputs x from Alice and y from Bob, they output values a and b such that the resulting probability distribution p(a, b|x, y) cannot be reproduced classically without communication, yet cannot itself be used to communicate (see Figure 2).

The original motivation for studying nonlocality is quantum mechanics. Indeed, in quantum mechanics, entanglement allows us to achieve nonlocal correlations similar to those described above. John Bell [3] was the first

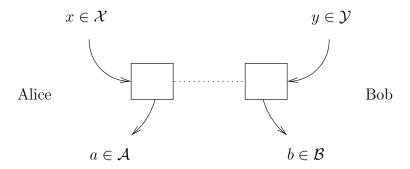


Figure 2: A generic nonlocal box. Alice and Bob input x and y respectively, and receive a and b respectively. The resulting probability distribution p(a, b|x, y) cannot be reproduced classically without communication, and yet does not itself allow Alice and Bob to communicate.

to show that measurement on shared entangled quantum state could produce correlations that cannot be locally simulated classically. Later, Clauser, Horne, Shimony and Holt [2] came up with an inequality (the so-called CHSH inequality) that provided a condition for a certain type of correlations to be explainable by classical means alone, and showed that quantum mechanics violated it in some cases. The mod2NLB was directly inspired by this inequality: the mod2NLB violates the CHSH inequality to its maximal algebraic value while quantum mechanics can be used to simulate a mod2NLB with up to approximately 85% efficiency [4].

The nonlocality of quantum mechanics has been known for a long time, but has only recently started to be studied by itself, i.e. independently from the study of entanglement. It is hoped that such an independent study will allow us to understand the implications of nonlocality in quantum mechanics more thoroughly. Furthermore, there is proof that entanglement and nonlocality are not the same. The first example came from [5], where it was proved that a single mod2NLB is not sufficient to simulate a non-maximally entangled pair of qubits, even though a perfect simulation of all correlations of the maximally entangled state of two qubits is possible with only one mod2NLB[6]. The final proof that entanglement and nonlocality are different resources came in [7], where it was proven that a simulation of n maximally entangled pair of qubits required $\Omega(2^n)$ mod2NLBs.

This asserts that entanglement and nonlocality should be treated as different types of resources. But while we know a fair bit about entanglement, comparatively speaking little is known about nonlocality. For instance, we have been able to isolate the maximally entangled state of two qubits as the "unit" of bipartite entanglement, since, together with local operations and classical communication, it allows us to create any other entangled state, provided we have enough copies. Is there an analogous concept for nonlocality? Would it be possible to identify a similar "unit of nonlocality", that would allow us to create other bipartite nonlocal correlations? The mod2NLB was the obvious candidate: its minimal size (binary inputs and outputs) and the fact that it violates the CHSH inequality, the only nontrivial inequality at these dimensions, maximally made it very attractive from that point of view.

There are more encouraging signs to support the mod2NLB's claim as the universal resource of nonlocality. One particularly interesting result is that the mod2NLB makes communication complexity [8] trivial [9, 10]. That is, if two players are allowed to use mod2NLBs, they can compute any boolean function of their inputs with a single bit of communication, regardless of what the function is.

In light of these facts, it is tempting to think that the mod2NLB could be considered as a unit of nonlocality that can be used to generate any other bipartite nonlocal correlation. Some progress has been made in this direction: in [11], Barrett and Pironio have shown that mod2NLBs alone can be used to simulate any two-output bipartite boxes. However, they have shown that there exist multipartite nonlocal correlations that cannot be simulated by mod2NLBs alone. What about the bipartite scenario? In [12], a family of bipartite nonlocal boxes is presented which can generate every two-input bipartite box. In this paper, we present a complementary negative result: we show that no *finite* set containing any general bipartite nonlocal boxes can simulate all bipartite nonlocal boxes.

We start, for intuition, by proving the non-universality of the traditional nonlocal box in Section 2. We prove that a finite number of mod2NLB cannot perfectly simulate the mod3NLB, to be defined at the beginning of Section 2. In Section 3, we generalize the result by proving that no finite set of finite-output-alphabet nonlocal boxes can be universal. We then conclude in Section 4.

2 The non-universality of the traditional nonlocal box

We will first begin by introducing the mod3NLB: $x \in \{0, 1\}, y \in \{0, 1\}, a \in \{0, 1, 2\}, b \in \{0, 1, 2\}, and$

$$p(a, b|x, y) = \begin{cases} \frac{1}{3} & \text{if } b - a = xy \mod 3\\ 0 & \text{otherwise} \end{cases}$$
 (1)

See Figure 3 for a graphical representation of the general modpNLB. Clearly, this does not allow communication between Alice and Bob, since, taken alone, a is completely independent from x and y, and likewise for b. The mod3NLB is therefore a valid nonlocal box and a simple extension of the traditional mod2NLB. It would seem reasonable, especially in light of [11], that such a nonlocal box could be simulated by mod2NLBs. However, the following theorem states the opposite.

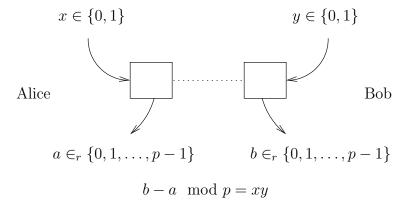


Figure 3: A graphical description of the mod pNLB.

Theorem 1. It is impossible to simulate the mod3NLB exactly using a finite number of mod2NLBs, infinite shared randomness and no communication between the two players.

Proof. Let's assume that there exists an algorithm (which may be probabilistic) that can perfectly simulate one instance of the mod3NLB using N mod2NLBs; we will then show that this assumption leads to a contradiction.

First, we can reduce the problem to deterministic algorithms in the following manner: any probabilistic algorithm can be represented as a collection of deterministic algorithms α_i , each with a certain probability of being selected. Since we require perfect simulation of the mod3NLB, the outputs of the algorithm must satisfy the equation b-a=xy with probability 1; hence each algorithm α_i with nonzero probability in any probabilistic algorithm must also satisfy this equation with probability 1. For our contradiction, we can therefore restrict ourselves to deterministic algorithms, since a correct probabilistic algorithm exists only if a deterministic algorithm satisfying b-a=xy exists.

Observe first that, for all deterministic algorithms, the output a is completely determined by x and the N output bits that Alice got from the mod2NLBs, since we can simulate Alice's algorithm using only those N+1 values. Likewise, we can do this on Bob's side to determine b from y and Bob's mod2NLB outputs. To formalize this, let z_A be the bit-string that Alice obtained from the mod2NLBs, and z_B be Bob's bit-string. Then there exist two functions F_A and F_B such that $a = F_A(x, z_A)$ and $b = F_B(y, z_B)$. Note that z_A and z_B are uniformly distributed on $\{0,1\}^N$. We can now define the following two probability distributions:

$$p_A(a|x) = \Pr\{F_A(x,Z) = a\}$$
 (2)

$$p_B(b|y) = \Pr\{F_B(y,Z) = b\}$$
(3)

where Z is a random variable uniformly distributed on $\{0,1\}^N$.

Let us note that 2^N is not divisible by 3, therefore p_A and p_B cannot be uniform for any value of x and y. Since we must be able to simulate the box perfectly, we must at least have:

$$p_A(q|0) = p_B(q|0) \tag{4}$$

$$p_A(q|0) = p_B(q|1) \tag{5}$$

$$p_A(q|1) = p_B(q|0) (6)$$

$$p_A(q|1) = p_B(q+1|1)$$
 (7)

where additions are performed mod 3. Condition (4) comes from the fact that if x = y = 0, then b = a every time, hence the two marginal distributions must be identical, and therefore $p_A(q|0) = p_B(q|0)$. The other three conditions correspond to similar conditions when the inputs are (a, b) = (0, 1), (a, b) = (1, 0) and (a, b) = (1, 1) respectively.

These conditions lead to a contradiction: (4) and (6) imply that $p_A(q|0) = p_A(q|1)$, which means that (5) and (7) imply that $p_B(q|1) = p_B(q+1|1)$. Since

 $p_B(q|1)$ cannot be uniform, we are forced to conclude that perfect simulation of the mod3NLB with N mod2NLBs is impossible.

3 Generalization to a finite set of nonlocal boxes

The result of Section 2 can be generalized to a finite set of nonlocal boxes, as defined in Section 1 and represented in Figure 2, where the dimensions of the output sets are finite, i.e. $|\mathcal{A}|, |\mathcal{B}| < \infty$. Before turning to the main theorem and its proof, we need to define the mod pNLB in the following manner:

$$p(a, b|x, y) = \begin{cases} \frac{1}{p} & \text{if } b - a = xy \mod p \\ 0 & \text{otherwise} \end{cases} , \tag{8}$$

where $|\mathcal{X}| = |\mathcal{Y}| = 2$ and $|\mathcal{A}| = |\mathcal{B}| = p$. This family of nonlocal boxes was first defined in [12]. It was also shown that this family, which is an infinite set, could be used to simulate any two-input bipartite nonlocal box. They also showed that a mod pNLB and a mod pNLB could simulate a mod pNLB, where r = pq. Here, we shall prove that for any finite set of nonlocal boxes, whatever their nature, there exist a member of this family that cannot be simulated by a finite number of boxes from the set. Before turning to our main theorem and its proof, we first need two technical lemmas.

Lemma 2. If a first-order formula ϕ is true in \mathbb{C} , it is then true in another field \mathfrak{F}_q which has a large enough characteristic q.

Proof. Recall that the characteristic k of a field \mathfrak{F} is the smallest positive integer such that for every $x \in \mathfrak{F}$, $x+\cdots+x=0$ where x is summed k times; if no such k exists, the characteristic of \mathfrak{F} is zero by definition. It is a known result in mathematical logic that the theory of algebraically closed fields with specified characteristic is complete (see, e.g., [13], pp. 178–179). This means that given a set of axioms which define an algebraically closed field of characteristic k, then every first-order formula can either be proven true or false from these axioms alone. Thus, since $\mathbb C$ is an algebraically closed field of characteristic 0, if a first-order formula is true in $\mathbb C$, it must be provable in the theory of algebraically closed fields of characteristic 0.

Now, the axioms of the theory of an algebraically closed field of characteristic 0 consist of two disjoint sets:

- A set of axioms S_{alb} which define an algebraically closed field
- A set of axioms $S_{\text{chr}} = \{\tau_p : p \text{ prime}\}$ where τ_p is the axiom that the field is not of characteristic p.

The proof of ϕ in this theory is finite, and thus uses a finite subset of axioms S. Specifically it only uses a finite subset of S_{chr} . Let q be a prime large enough such that $\tau_p \notin S$ for every $p \geq q$, and let \mathfrak{F}_q be any algebraically closed field of characteristic q. The proof only uses axioms which hold in \mathfrak{F}_q and, therefore, the statement ϕ is true in \mathfrak{F}_q .

Lemma 3. No finitely generated extension ring of \mathbb{Z} contains all numbers of the form 1/p, where p is a prime number.

Proof. Let T be some extension ring of \mathbb{Z} generated by a finite set $G = \{q_1, \ldots, q_n\} \subset \mathbb{R}$ over \mathbb{Z} , meaning that T is composed of numbers which are the sum of products of numbers in G and \mathbb{Z} . Assume, by contradiction, that all numbers of the form 1/p, where p is prime, are contained in T. Let \mathfrak{I} be the set of all n-variable polynomial relations with coefficients in \mathbb{Z} satisfied by q_1, \ldots, q_n ; it can be shown that \mathfrak{I} is an ideal. From the Hilbert basis theorem [14], it follows that \mathfrak{I} is finitely generated; this means that there exists a set of polynomials $F = \{f_1, \ldots, f_m\}$ such that every P in \mathfrak{I} can be written as $P = \sum_{k=1}^m r_k f_k$, where $\{r_1, \ldots, r_m\}$ is a set of polynomials that depends on P.

Now, consider the following first-order formula;

$$\phi: \exists y_1, \dots, y_n [f_1(y_1, \dots, y_n) = \dots = f_m(y_1, \dots, y_n) = 0].$$
 (9)

Obviously, ϕ must be true in \mathbb{R} , since q_1, \ldots, q_n satisfies the formula. It must therefore also be true in \mathbb{C} and, by Lemma 2, in a field \mathfrak{F}_p of sufficiently large characteristic p.

Let y_1, \ldots, y_n be the elements of \mathfrak{F}_p guaranteed by ϕ , which satisfy f_1, \ldots, f_m . By hypothesis, $1/p \in T$, so 1/p can be expressed as a sum of products of numbers in G. Thus, there is some n-variable polynomial g with integer coefficients such that $g(q_1, \ldots, q_n) = 1/p$ and therefore $1 - pg(q_1, \ldots, q_n) = 0$. This means that $1 - pg \in \mathfrak{I}$. We then have that $1 - pg(y_1, \ldots, y_n) = 0$ in \mathfrak{F}_p . But in \mathfrak{F}_p , p = 0 by definition and $1 - pg(y_1, \ldots, y_n) = 1$. Thus we have a contradiction.

Theorem 4. Let S be a finite set of nonlocal boxes, each with a finite output alphabets A and B. Then there exists p such that the modpNLB cannot be simulated by a finite number of nonlocal boxes taken from the set S.

Proof. Let us first note that, as in the Proof of Theorem 1, the fact that the relation of the modpNLB must be perfectly simulated entails that we can consider only deterministic protocols; meaning that the only source of randomness is the nonlocal boxes from the set S, that we do not have any shared random variables and that the functions F_A and F_B , with which we calculate the final output of our simulation from the internal nonlocal boxes and the input, are deterministic. Let $p_k^{(l)}$ be the marginal probability of Alice's output k of the nonlocal box l from the set S. Likewise, let $q_k^{(l)}$ be the marginal probability of Bob's output k of the nonlocal box l from the set S. Let $P = \{p_1^{(1)}, p_2^{(1)}, \ldots, p_1^{(2)}, \ldots\}$ and $Q = \{q_1^{(1)}, q_2^{(1)}, \ldots, q_1^{(2)}, \ldots\}$ be the total collection of probabilities for outputs of S. P and Q are obviously finite, since each output alphabet and S are finite.

Let a be some element in the modpNLB output alphabet. From the fact that the output of the simulation is calculated from a deterministic function on the input and the output of the internal nonlocal boxes, the marginal probability of a must be some linear combination (with integral coefficients) of products of probabilities in the original set. Mathematically, a must be an element of the integral extension ring generated by P over \mathbb{Z} . In other words, a is in the "set" of numbers generated the addition of products of numbers from P and \mathbb{Z} .

However, it is proven in Lemma 3 that \mathbb{Q} is not a finitely generated extension ring of \mathbb{Z} . Meaning that it is impossible to create every number in \mathbb{Q} by the addition of products of numbers from P and \mathbb{Z} . Therefore, there exists some $q = 1/p \in \mathbb{Q}$ which is not in the integral extension ring generated by P over \mathbb{Z} . Hence, the modpNLB cannot be simulated by a finite number of nonlocal boxes from the set S.

4 Discussion and Conclusion

We have proven that no finite set of finite-output-alphabet nonlocal boxes can be universal. Therefore only nonlocal boxes with an infinite number of outputs can be. However, there exists no compelling candidates, and one might argue that such a box would be even more artificial and less elegant than the traditional mod2NLB.

It is to be noted that our result does not contradict those of [11], since the universality of the family of mod pNLBs is defined for binary input nonlocal boxes and requires an infinite set of mod pNLBs.

Our result exhibits yet a new difference between nonlocality and entanglement. The latter has a very simple and attractive universal resource, the maximally entangled pair of qubits, while the former has no such things. Our result also suggest that one must be careful about statements made with the traditional mod2NLB, for it cannot be associated with a general idea of nonlocality. It is important to stress at this point that we do not believe research in nonlocal boxes to be futile. For example, one can still uncover some intuitions about Nature when studying the mod2NLB. In [15], it was proven, using mod2NLBs, that if quantum mechanics were slightly more nonlocal, it would have drastic and arguably unbelievable consequences in communication complexity.

This work raises a philosophical question. What is the difference between entanglement and nonlocality? Why does entanglement have a universal resource while nonlocality doesn't? It is tempting to think that it might be related to the fact that we limited the output dimensions of our nonlocal boxes while measurements on entangled states can have any number of outputs. However, we would like to point out that the quantum universality theorem uses quantum teleportation as its main building block, which requires measurements with a finite set of possible outputs. We believe that the answer might be related to the question of the difference between entanglement measures and nonlocality measures [16]. In our scenario, we do not allow the participants to use classical communication, since it is a nonlocal resource. On the other hand, the universality of the maximally entangled pair of qubits is established by allowing the participants any resource that do not increase entanglement: shared randomness, local operations and classical communication. If we take away that last resource, the universality theorem of entanglement breaks down. Therefore, nonlocality is directly used in order to generate any possible entangled state out of a maximally entangled pair of qubits. What does this entail precisely? We will let the reader ponder this question.

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References

- [1] S. Popescu and D. Rohrlich, "Quantum nonlocality as an axiom," Foundations of Physics, vol. 24, pp. 379–385, 1994.
- [2] J. F. Clauser, M. A. Horne, A. Shimony, and R. A. Holt, "Proposed experiment to test local hidden-variable theories," *Physical Review Letters*, vol. 23, pp. 880–884, 1969.
- [3] J. Bell, "On the Einstein-Podolsky-Rosen paradox," *Physics*, vol. 1, pp. 195–200, 1964.
- [4] B. S. Cirel'son, "Quantum generalizations of Bell's inequality," *Letters in Mathematical Physics*, vol. 4, pp. 93–100, 1980.
- [5] N. Brunner, N. Gisin, and V. Scarani, "Entanglement and non-locality are different resources," *New Journal of Physics*, vol. 7, p. 88, 2005.
- [6] N. Cerf, N. Gisin, S. Massar, and S. Popescu, "Quantum entanglement can be simulated without communication," *Physical Review Letters*, vol. 94, p. 220403, 2005.
- [7] A. Broadbent and A. A. Méthot, "On the power of non-local boxes," *Theoretical Computer Science C*, vol. 358, pp. 3–14, 2006.
- [8] E. Kushilevitz and N. Nisan, *Communication Complexity*. Cambridge: Cambridge University Press, 1997.
- [9] W. van Dam, Nonlocality & Communication Complexity. Oxford University: Ph.D. Thesis, 2000.
- [10] R. Cleve, personal communication.

- [11] J. Barrett and S. Pironio, "Popescu-Rohrlich correlations as a unit of nonlocality," *Physical Review Letters*, vol. 95, p. 140401, 2005.
- [12] J. Barrett, N. Linden, S. Massar, S. Pironio, S. Popescu, and D. Roberts, "Non-local correlations as an information theoretic resource," *Physical Review A*, vol. 71, p. 022101, 2005.
- [13] J. L. Bell and A. B. Slomson, *Models and ultraproducts*. North-Holland, Amsterdam, 1969.
- [14] D. A. Cox, J. B. Little, and D. O'Shea, *Ideals, Varieties, and Algorithms*. Springer-Verlag, 1997.
- [15] G. Brassard, H. Buhrman, N. Linden, A. A. Méthot, A. Tapp, and F. Unger, "A limit on non-local correlations in any world where communication complexity is not trivial," *Physical Review Letters*, vol. 96, p. 205401, 2006, reprinted in *Virtual Journal of Quantum Information*, vol. 6, 2006.
- [16] A. A. Méthot and V. Scarani, "An anomaly of non-locality," Quantum Information and Computation, vol. 7, pp. 157–170, 2007.